

P200S - January 15, 2018

Item # P200S was discontinued on January 15, 2018. For informational purposes, this is a copy of the website content at that time and is valid only for the stated product.

PRECISION PINHOLES AND PINHOLE WHEEL

- ▶ Standard and High-Power Precision Pinholes Mounted in Ø1" Disks
- ▶ Single Pinhole Sizes from Ø1 µm to Ø1 mm
- ▶ 16-Position Pinhole Wheel with Hole Sizes from Ø25 µm to Ø2 mm



OVERVIEW

Features

- Mounted Precision Pinholes from Ø1 µm to Ø1 mm
- Mounted High-Power Pinholes from Ø10 µm to Ø50 µm
- Chrome-Plated Fused Silica Pinhole Wheel with 16 Pinholes from Ø25 µm to Ø2 mm

Single Precision Pinholes

Single mounted precision pinholes are available with pinhole diameters from 1 µm to 1 mm. We also offer high-power versions with pinhole diameters from 10 µm to 50 µm. For many applications, such as holography, spatial intensity variations in the laser beam are unacceptable. Using precision pinholes in conjunction with positioning and focusing equipment such as our KT310(/M) Spatial Filter System creates a "noise" filter, effectively stripping variations in intensity out of a Gaussian beam. Please see the *Tutorial* tab for more information on spatial filters.

If you do not see what you need in our stocked offerings below, it is possible to special order pinholes that are fabricated from different substrate materials, have different pinhole sizes, incorporate multiple holes in one foil, or provide different pinhole configurations. Customized pinhole housings are also available. Please contact Tech Support to discuss your specific needs.

Pinhole Wheels

In addition to single pinholes, Thorlabs offers pinhole wheels that contain 16 radially-spaced pinholes that are lithographically etched onto a chrome-plated fused silica substrate. These wheels allow the user to test multiple pinhole sizes within a setup. The 16 pinholes range in size from Ø25 µm to Ø2 mm and both sides of the wheel are AR coated for 350 - 700 nm.

TUTORIAL

Principles of Spatial Filters

For many applications, such as holography, spatial intensity variations in the laser beam are unacceptable. Our KT310 spatial filter system is ideal for producing a clean Gaussian beam.

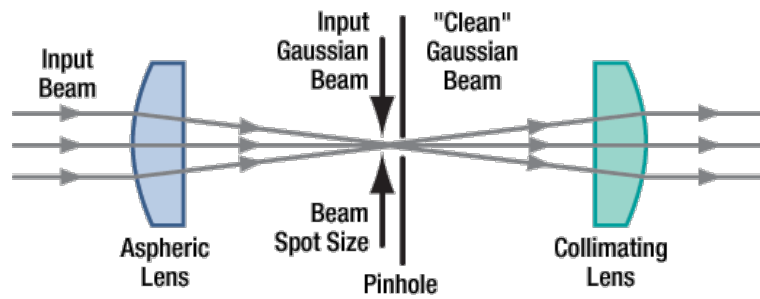


Figure 1: Spatial Filter System

The input Gaussian beam has spatially varying intensity "noise". When a beam is focused by an aspheric lens, the input beam is transformed into a central Gaussian spot (on the optical axis) and side fringes, which represent the unwanted "noise" (see Figure 2 below). The radial position of the side fringes is proportional to the spatial frequency of the "noise".

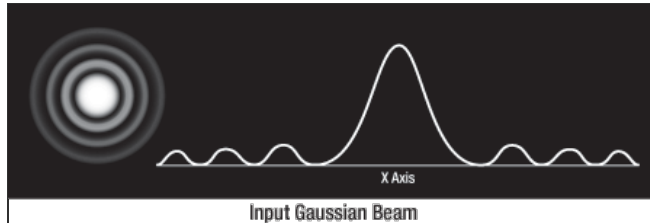


Figure 2

By centering a pinhole on a central Gaussian spot, the "clean" portion of the beam can pass while the "noise" fringes are blocked (see Figure 3 below).

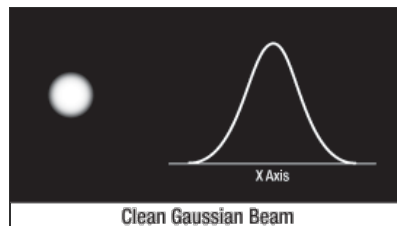


Figure 3

The diffraction-limited spot size at the 99% contour is given by:

$$D = \frac{\lambda f}{r}$$

where λ = wavelength, f =focal length and r = input beam radius at the $1/e^2$ point.

Choosing the Correct Optics and Pinhole for Your Spatial Filter System

The correct optics and pinhole for your application depend on the input wavelength, source beam diameter, and desired exit beam diameter.

For example, suppose that you are using a 650 nm diode laser source that has a diameter ($1/e^2$) of 1.2 mm and want your beam exiting the spatial filter system to be about 4.4 mm in diameter. Based on these parameters, the C560TME-B mounted aspheric lens would be an appropriate choice for the input side of spatial filter system because it is designed for use at 650 nm, and its clear aperture measures 5.1 mm, which is large enough to accommodate the entire diameter of the laser source.

The equation for diffraction limited spot size at the 99% contour is given above, and for this example, $\lambda = (650 \times 10^{-9} \text{ m})$, $f = 13.86 \text{ mm}$ for the C560TM-B, and $r = 0.6 \text{ mm}$. Substitution yields

$$D = \frac{(650 \times 10^{-9} \text{ m})(13.86 \text{ mm})}{0.6 \text{ mm}} \approx 15 \mu\text{m}$$

Diffraction-Limited Spot Size (650 nm source, Ø1.2 mm beam)

The pinhole should be chosen so that it is approximately 30% larger than D . If the pinhole is too small, the beam will be clipped, but if it is too large, more than the TEM_{00} mode will get through the pinhole. Therefore, for this example, the pinhole should ideally be 19.5 microns. Hence, we would recommend the mounted pinhole P20H, which has a pinhole size of 20 μm . Parameters that can be changed to alter the beam waist diameter, and thus the pinhole size required, include changing the input beam diameter and focal length of focusing lens. Decreasing the input beam diameter will increase the beam waist

diameter. Using a longer focal length focusing lens will also increase the beam waist diameter.

Finally, we need to choose the optic on the output side of the spatial filter so that the collimated beam's diameter is the desired 4.4 mm. To determine the correct focal length for the lens, consider the following diagram in Figure 4, which is not drawn to scale. From the triangle on the left-hand side, the angle is determined to be approximately 2.48° . Using this same angle for the triangle on the right-hand side, the focal length for the plano-convex lens should be approximately 50 mm.

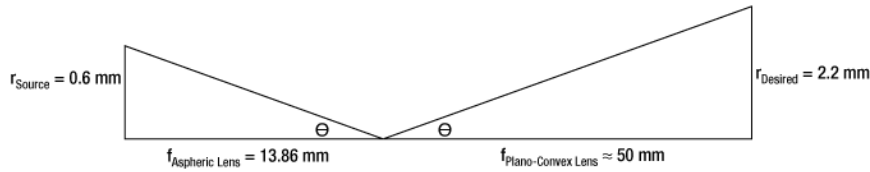


Figure 4: Beam Expansion Example

For this focal length, we recommend the LA1131-B plano-convex lens [with $f = 50$ mm at the design wavelength ($\lambda = 633$ nm), this is still a good approximation for f at the source wavelength ($\lambda = 650$ nm)].

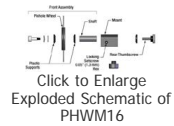
Note: The beam expansion equals the focal length of the output side divided by the focal length of the input side.

For optimal performance, a large-diameter aspheric lens can be used in place of a plano-convex lens if the necessary focal length on the output side is 20 mm (see AL2520-A, AL2520-B, AL2520-C). These lenses are 25 mm in diameter and can be held in place using the supplied SM1RR Retaining Ring.

ASSEMBLY

Mounting the Pinhole Wheel

The 16-position pinhole wheel can be post mounted using the NDC-PM Post Mount Assembly (included with the PHWM16) by following the assembly steps below. Adapters for both 8-32 and M4 mounting holes are included.



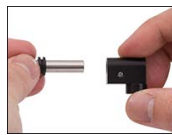
1. Tighten the 4-40 locking setscrew on the side of the NDC-PM assembly using the included 0.05" (1.3 mm) hex key to lock the shaft in position. Unscrew and remove the rear thumbscrew, then loosen the locking setscrew (Figure 1).
2. Pull the front assembly from the mount (Figure 2).
3. While holding the assembly vertically unscrew the shaft. Remove the top plastic support and place the pinhole wheel with the engraving facing down. Then, secure the pinhole wheel by placing the plastic support in position and screwing the shaft onto the cap screw (Figure 3).
4. Insert the front assembly with pinhole wheel into the mount, lock the locking setscrew, and then screw on the rear thumbscrew (Figure 4).

During installation and use, the guidelines below may be helpful:

- Holding the mount vertically (see the photo below) prevents the components (e.g., washer, spacer, and plastic supports) from falling apart during assembly.
- Use light force when securing the pinhole wheel as overtightening may cause the wheel to crack.
- The wheel may still rotate if moved by hand even if the locking setscrew is fully tightened.



Click to Enlarge
Figure 1: Lock the Mount and Remove the Rear Thumbscrew



Click to Enlarge
Figure 2: Unlock the Mount and Pull the Front Assembly Out of the Mount



Click to Enlarge
Figure 3: Hold Vertically and Secure the Pinhole Wheel



Click to Enlarge
Figure 4: Insert the Front Assembly into the Mount and Screw on the Rear Thumbscrew

DAMAGE THRESHOLDS

Damage Threshold Data for Thorlabs' Pinholes and Pinhole Wheel

The specifications to the right are measured data for Thorlabs' high-power single pinholes and 16-position pinhole wheel.

Damage Threshold Specifications	
Item #	Damage Threshold
P10CH, P25CH, P50CH	5×10^5 W/mm ² , 75 ns Pulse @ 700 nm
	1×10^6 W/mm ² , 10 ns Pulse @ 700 nm
	10 W/mm ² , CW @ 10.6 μ m
PHW16, PHWM16	7.5 J/cm ² at 532 nm, 10 ns, 10 Hz, \varnothing 0.491 mm

Laser Induced Damage Threshold Tutorial

The following is a general overview of how laser induced damage thresholds are measured and how the values may be utilized in determining the appropriateness of an optic for a given application. When choosing optics, it is important to understand the Laser Induced Damage Threshold (LIDT) of the optics being used. The LIDT for an optic greatly depends on the type of laser you are using. Continuous wave (CW) lasers typically cause damage from thermal effects (absorption either in the coating or in the substrate). Pulsed lasers, on the other hand, often strip electrons from the lattice structure of an optic before causing thermal damage. Note that the guideline presented here assumes room temperature operation and optics in new condition (i.e., within scratch-dig spec, surface free of contamination, etc.). Because dust or other particles on the surface of an optic can cause damage at lower thresholds, we recommend keeping surfaces clean and free of debris. For more information on cleaning optics, please see our *Optics Cleaning* tutorial.

Testing Method

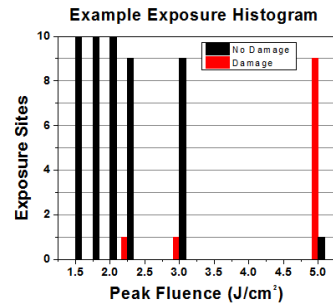
Thorlabs' LIDT testing is done in compliance with ISO/DIS11254 and ISO 21254 specifications.

First, a low-power/energy beam is directed to the optic under test. The optic is exposed in 10 locations to this laser beam for 30 seconds (CW) or for a number of pulses (pulse repetition frequency specified). After exposure, the optic is examined by a microscope (~100X magnification) for any visible damage. The number of locations that are damaged at a particular power/energy level is recorded. Next, the power/energy is either increased or decreased and the optic is exposed at 10 new locations. This process is repeated until damage is observed. The damage threshold is then assigned to be the highest power/energy that the optic can withstand without causing damage. A histogram such as that below represents the testing of one BB1-E02 mirror.



The photograph above is a protected aluminum-coated mirror after LIDT testing. In this particular test, it handled 0.43 J/cm² (1064 nm, 10 ns pulse, 10 Hz, Ø1.000 mm) before damage.

According to the test, the damage threshold of the mirror was 2.00 J/cm² (532 nm, 10 ns pulse, 10 Hz, Ø0.803 mm). Please keep in mind that these tests are performed on clean optics, as dirt and contamination can significantly lower the damage threshold of a component. While the test results are only representative of one coating run, Thorlabs specifies damage threshold values that account for coating variances.



Example Test Data			
Fluence	# of Tested Locations	Locations with Damage	Locations Without Damage
1.50 J/cm ²	10	0	10
1.75 J/cm ²	10	0	10
2.00 J/cm ²	10	0	10
2.25 J/cm ²	10	1	9
3.00 J/cm ²	10	1	9
5.00 J/cm ²	10	9	1

Continuous Wave and Long-Pulse Lasers

When an optic is damaged by a continuous wave (CW) laser, it is usually due to the melting of the surface as a result of absorbing the laser's energy or damage to the optical coating (antireflection) [1]. Pulsed lasers with pulse lengths longer than 1 µs can be treated as CW lasers for LIDT discussions.

When pulse lengths are between 1 ns and 1 µs, laser-induced damage can occur either because of absorption or a dielectric breakdown (therefore, a user must check both CW and pulsed LIDT). Absorption is either due to an intrinsic property of the optic or due to surface irregularities; thus LIDT values are only valid for optics meeting or exceeding the surface quality specifications given by a manufacturer. While many optics can handle high power CW lasers, cemented (e.g., achromatic doublets) or highly absorptive (e.g., ND filters) optics tend to have lower CW damage thresholds. These lower thresholds are due to absorption or scattering in the cement or metal coating.

Pulsed lasers with high pulse repetition frequencies (PRF) may behave similarly to CW beams. Unfortunately, this is highly dependent on factors such as absorption and thermal diffusivity, so there is no reliable method for determining when a high PRF laser will damage an optic due to thermal effects. For beams with a high PRF both the average and peak powers must be compared to the equivalent CW power. Additionally, for highly transparent materials, there is little to no drop in the LIDT with increasing PRF.

In order to use the specified CW damage threshold of an optic, it is necessary to know the following:

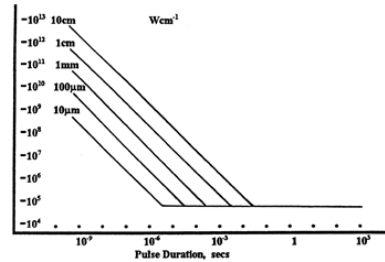
1. Wavelength of your laser
2. Beam diameter of your beam (1/e²)
3. Approximate intensity profile of your beam (e.g., Gaussian)
4. Linear power density of your beam (total power divided by 1/e² beam diameter)

Thorlabs expresses LIDT for CW lasers as a linear power density measured in W/cm. In this regime, the LIDT given as a linear power density can be applied to any beam diameter; one does not need to compute an adjusted LIDT to adjust for changes in spot size, as demonstrated by the graph to the right. Average linear power density can be calculated using the equation below.

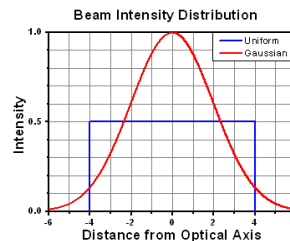
$$\text{Linear Power Density} = \frac{\text{Power}}{\text{Beam Diameter}}$$

The calculation above assumes a uniform beam intensity profile. You must now consider hotspots in the beam or other non-uniform intensity profiles and roughly calculate a maximum power density. For reference, a Gaussian beam typically has a maximum power density that is twice that of the uniform beam (see lower right).

Now compare the maximum power density to that which is specified as the LIDT for the optic. If the optic was tested at a wavelength other than your operating wavelength, the damage threshold must be scaled appropriately. A good rule of thumb is that the damage threshold has a linear relationship with wavelength such that as you move to shorter wavelengths, the damage threshold decreases (i.e., a LIDT of 10 W/cm at 1310 nm scales to 5 W/cm at 655 nm):



LIDT in linear power density vs. pulse length and spot size. For long pulses to CW, linear power density becomes a constant with spot size. This graph was obtained from [1].



$$\text{Adjusted LIDT} = \text{LIDT Power} \left(\frac{\text{Your Wavelength}}{\text{LIDT Wavelength}} \right)$$

While this rule of thumb provides a general trend, it is not a quantitative analysis of LIDT vs wavelength. In CW applications, for instance, damage scales more strongly with absorption in the coating and substrate, which does not necessarily scale well with wavelength. While the above procedure provides a good rule of thumb for LIDT values, please contact Tech Support if your wavelength is different from the specified LIDT wavelength. If your power density is less than the adjusted LIDT of the optic, then the optic should work for your application.

Please note that we have a buffer built in between the specified damage thresholds online and the tests which we have done, which accommodates variation between batches. Upon request, we can provide individual test information and a testing certificate. The damage analysis will be carried out on a similar optic (customer's optic will not be damaged). Testing may result in additional costs or lead times. Contact Tech Support for more information.

Pulsed Lasers

As previously stated, pulsed lasers typically induce a different type of damage to the optic than CW lasers. Pulsed lasers often do not heat the optic enough to damage it; instead, pulsed lasers produce strong electric fields capable of inducing dielectric breakdown in the material. Unfortunately, it can be very difficult to compare the LIDT specification of an optic to your laser. There are multiple regimes in which a pulsed laser can damage an optic and this is based on the laser's pulse length. The highlighted columns in the table below outline the relevant pulse lengths for our specified LIDT values.

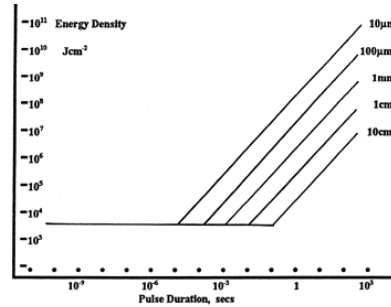
Pulses shorter than 10^{-9} s cannot be compared to our specified LIDT values with much reliability. In this ultra-short-pulse regime various mechanics, such as multiphoton-avalanche ionization, take over as the predominate damage mechanism [2]. In contrast, pulses between 10^{-7} s and 10^{-4} s may cause damage to an optic either because of dielectric breakdown or thermal effects. This means that both CW and pulsed damage thresholds must be compared to the laser beam to determine whether the optic is suitable for your application.

Pulse Duration	$t < 10^{-9}$ s	$10^{-9} < t < 10^{-7}$ s	$10^{-7} < t < 10^{-4}$ s	$t > 10^{-4}$ s
Damage Mechanism	Avalanche Ionization	Dielectric Breakdown	Dielectric Breakdown or Thermal	Thermal
Relevant Damage Specification	No Comparison (See Above)	Pulsed	Pulsed and CW	CW

When comparing an LIDT specified for a pulsed laser to your laser, it is essential to know the following:

1. Wavelength of your laser
2. Energy density of your beam (total energy divided by $1/e^2$ area)
3. Pulse length of your laser
4. Pulse repetition frequency (prf) of your laser
5. Beam diameter of your laser ($1/e^2$)
6. Approximate intensity profile of your beam (e.g., Gaussian)

The energy density of your beam should be calculated in terms of J/cm^2 . The graph to the right shows why expressing the LIDT as an energy density provides the best metric for short pulse sources. In this regime, the LIDT given as an energy density can be applied to any beam diameter; one does not need to compute an adjusted LIDT to adjust for changes in spot size. This calculation assumes a uniform beam intensity profile. You must now adjust this energy density to account for hotspots or other nonuniform intensity profiles and roughly calculate a maximum energy density. For reference a Gaussian beam typically has a maximum energy density that is twice that of the $1/e^2$ beam.



LIDT in energy density vs. pulse length and spot size. For short pulses, energy density becomes a constant with spot size. This graph was obtained from [1].

Now compare the maximum energy density to that which is specified as the LIDT for the optic. If the optic was tested at a wavelength other than your operating wavelength, the damage threshold must be scaled appropriately [3]. A good rule of thumb is that the damage threshold has an inverse square root relationship with wavelength such that as you move to shorter wavelengths, the damage threshold decreases (i.e., a LIDT of $1 J/cm^2$ at 1064 nm scales to $0.7 J/cm^2$ at 532 nm):

$$\text{Adjusted LIDT} = \text{LIDT Energy} \sqrt{\frac{\text{Your Wavelength}}{\text{LIDT Wavelength}}}$$

You now have a wavelength-adjusted energy density, which you will use in the following step.

Beam diameter is also important to know when comparing damage thresholds. While the LIDT, when expressed in units of J/cm^2 , scales independently of spot size; large beam sizes are more likely to illuminate a larger number of defects which can lead to greater variances in the LIDT [4]. For data presented here, a <1 mm beam size was used to measure the LIDT. For beams sizes greater than 5 mm, the LIDT (J/cm^2) will not scale independently of beam diameter due to the larger size beam exposing more defects.

The pulse length must now be compensated for. The longer the pulse duration, the more energy the optic can handle. For pulse widths between 1 - 100 ns, an approximation is as follows:

$$\text{Adjusted LIDT} = \text{LIDT Energy} \sqrt{\frac{\text{Your Pulse Length}}{\text{LIDT Pulse Length}}}$$

Use this formula to calculate the Adjusted LIDT for an optic based on your pulse length. If your maximum energy density is less than this adjusted LIDT maximum energy density, then the optic should be suitable for your application. Keep in mind that this calculation is only used for pulses between 10^{-9} s and 10^{-7} s. For pulses between 10^{-7} s and 10^{-4} s, the CW LIDT must also be checked before deeming the optic appropriate for your application.

Please note that we have a buffer built in between the specified damage thresholds online and the tests which we have done, which accommodates variation between batches. Upon request, we can provide individual test information and a testing certificate. Contact Tech Support for more information.

- [1] R. M. Wood, *Optics and Laser Tech.* **29**, 517 (1998).
 [2] Roger M. Wood, *Laser-Induced Damage of Optical Materials* (Institute of Physics Publishing, Philadelphia, PA, 2003).
 [3] C. W. Carr *et al.*, *Phys. Rev. Lett.* **91**, 127402 (2003).
 [4] N. Bloembergen, *Appl. Opt.* **12**, 661 (1973).

LIDT CALCULATIONS

In order to illustrate the process of determining whether a given laser system will damage an optic, a number of example calculations of laser induced damage threshold are given below. For assistance with performing similar calculations, we provide a spreadsheet calculator that can be downloaded by clicking the button to the right. To use the calculator, enter the specified LIDT value of the optic under consideration and the relevant parameters of your laser system in the green boxes. The spreadsheet will then calculate a linear power density for CW and pulsed systems, as well as an energy density value for pulsed systems. These values are used to calculate adjusted, scaled LIDT values for the optics based on accepted scaling laws. This calculator assumes a Gaussian beam profile, so a correction factor must be introduced for other beam shapes (uniform, etc.). The LIDT scaling laws are determined from empirical relationships; their accuracy is not guaranteed. Remember that absorption by optics or coatings can significantly reduce LIDT in some spectral regions. These LIDT values are not valid for ultrashort pulses less than one nanosecond in duration.

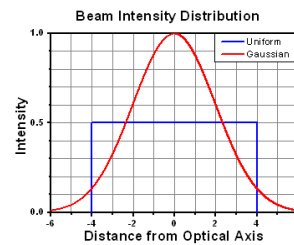
[LIDT Calculator](#)

CW Laser Example

Suppose that a CW laser system at 1319 nm produces a 0.5 W Gaussian beam that has a $1/e^2$ diameter of 10 mm. A naive calculation of the average linear power density of this beam would yield a value of 0.5 W/cm, given by the total power divided by the beam diameter:

$$\text{Linear Power Density} = \frac{\text{Power}}{\text{Beam Diameter}}$$

However, the maximum power density of a Gaussian beam is about twice the maximum power density of a uniform beam, as shown in the graph to the right. Therefore, a more accurate determination of the maximum linear power density of the system is 1 W/cm.



A Gaussian beam profile has about twice the maximum intensity of a uniform beam profile.

An AC127-030-C achromatic doublet lens has a specified CW LIDT of 350 W/cm, as tested at 1550 nm. CW damage threshold values typically scale directly with the wavelength of the laser source, so this yields an adjusted LIDT value:

$$\text{Adjusted LIDT} = \text{LIDT Power} \left(\frac{\text{Your Wavelength}}{\text{LIDT Wavelength}} \right)$$

The adjusted LIDT value of 350 W/cm \times (1319 nm / 1550 nm) = 298 W/cm is significantly higher than the calculated maximum linear power density of the laser system, so it would be safe to use this doublet lens for this application.

Pulsed Nanosecond Laser Example: Scaling for Different Pulse Durations

Suppose that a pulsed Nd:YAG laser system is frequency tripled to produce a 10 Hz output, consisting of 2 ns output pulses at 355 nm, each with 1 J of energy, in a Gaussian beam with a 1.9 cm beam diameter ($1/e^2$). The average energy density of each pulse is found by dividing the pulse energy by the beam area:

$$\text{Energy Density} = \frac{\text{Pulse Energy}}{\text{Beam Area}}$$

As described above, the maximum energy density of a Gaussian beam is about twice the average energy density. So, the maximum energy density of this beam is $\sim 0.7 \text{ J/cm}^2$.

The energy density of the beam can be compared to the LIDT values of 1 J/cm^2 and 3.5 J/cm^2 for a BB1-E01 broadband dielectric mirror and an NB1-K08 Nd:YAG laser line mirror, respectively. Both of these LIDT values, while measured at 355 nm, were determined with a 10 ns pulsed laser at 10 Hz. Therefore, an adjustment must be applied for the shorter pulse duration of the system under consideration. As described on the previous tab, LIDT values in the nanosecond pulse regime scale with the square root of the laser pulse duration:

$$\text{Adjusted LIDT} = \text{LIDT Energy} \sqrt{\frac{\text{Your Pulse Length}}{\text{LIDT Pulse Length}}}$$

This adjustment factor results in LIDT values of 0.45 J/cm^2 for the BB1-E01 broadband mirror and 1.6 J/cm^2 for the Nd:YAG laser line mirror, which are to be compared with the 0.7 J/cm^2 maximum energy density of the beam. While the broadband mirror would likely be damaged by the laser, the more specialized laser line mirror is appropriate for use with this system.

Pulsed Nanosecond Laser Example: Scaling for Different Wavelengths

Suppose that a pulsed laser system emits 10 ns pulses at 2.5 Hz, each with 100 mJ of energy at 1064 nm in a 16 mm diameter beam ($1/e^2$) that must be attenuated with a neutral density filter. For a Gaussian output, these specifications result in a maximum energy density of 0.1 J/cm^2 . The damage threshold of

an NDUV10A Ø25 mm, OD 1.0, reflective neutral density filter is 0.05 J/cm² for 10 ns pulses at 355 nm, while the damage threshold of the similar NE10A absorptive filter is 10 J/cm² for 10 ns pulses at 532 nm. As described on the previous tab, the LIDT value of an optic scales with the square root of the wavelength in the nanosecond pulse regime:

$$\text{Adjusted LIDT} = \text{LIDT Energy} \sqrt{\frac{\text{Your Wavelength}}{\text{LIDT Wavelength}}}$$

This scaling gives adjusted LIDT values of 0.08 J/cm² for the reflective filter and 14 J/cm² for the absorptive filter. In this case, the absorptive filter is the best choice in order to avoid optical damage.

Pulsed Microsecond Laser Example

Consider a laser system that produces 1 μs pulses, each containing 150 μJ of energy at a repetition rate of 50 kHz, resulting in a relatively high duty cycle of 5%. This system falls somewhere between the regimes of CW and pulsed laser induced damage, and could potentially damage an optic by mechanisms associated with either regime. As a result, both CW and pulsed LIDT values must be compared to the properties of the laser system to ensure safe operation.

If this relatively long-pulse laser emits a Gaussian 12.7 mm diameter beam (1/e²) at 980 nm, then the resulting output has a linear power density of 5.9 W/cm and an energy density of 1.2 × 10⁻⁴ J/cm² per pulse. This can be compared to the LIDT values for a WPQ10E-980 polymer zero-order quarter-wave plate, which are 5 W/cm for CW radiation at 810 nm and 5 J/cm² for a 10 ns pulse at 810 nm. As before, the CW LIDT of the optic scales linearly with the laser wavelength, resulting in an adjusted CW value of 6 W/cm at 980 nm. On the other hand, the pulsed LIDT scales with the square root of the laser wavelength and the square root of the pulse duration, resulting in an adjusted value of 55 J/cm² for a 1 μs pulse at 980 nm. The pulsed LIDT of the optic is significantly greater than the energy density of the laser pulse, so individual pulses will not damage the wave plate. However, the large average linear power density of the laser system may cause thermal damage to the optic, much like a high-power CW beam.

LAB FACTS

Comparison of Circularization Techniques for Elliptical Beams

Edge-emitting laser diodes emit elliptical beams as a consequence of the rectangular cross sections of their emission apertures. The component of the beam corresponding to the narrower dimension of the aperture has a greater divergence angle than the orthogonal beam component. As one component diverges more rapidly than the other, the beam shape is elliptical rather than circular.

Elliptical beam shapes can be undesirable, as the spot size of the focused beam is larger than if the beam were circular, and larger spot sizes have lower irradiances (power per area). Several different techniques can be used to circularize an elliptical beam, and we experimented with and compared the performance of three methods based on a pair of cylindrical lenses, an anamorphic prism pair, and a spatial filter. The characteristics of the circularized beams were evaluated by performing M² measurements, wavefront measurements, and measuring the transmitted power.

While we demonstrated that each circularization technique improves the circularity of the elliptical input beam, we showed that each technique provides a different balance of circularization, beam quality, and transmitted power. Our results, which are documented in this Lab Fact, indicate that an application's specific requirements will determine which is the best circularization technique to choose.

Experimental Design and Setup

The experimental setup is shown in the picture at the top-right. The elliptically-shaped, collimated beam of a temperature-stabilized 670 nm laser diode was input to each of our circularization systems. Collimation results in a low-divergence beam, but it does not affect the beam shape.

The beam circularization systems, shown to the right, were placed, one at a time, in the vacant spot in the setup highlighted by the yellow rectangle. With this arrangement, it was possible to use the same experimental conditions when evaluating each circularization technique, which allowed the performance of each to be directly compared with the others. Some information describing selection and configuration procedures for several components used in this experimental work can be accessed by clicking the following hyperlinks:

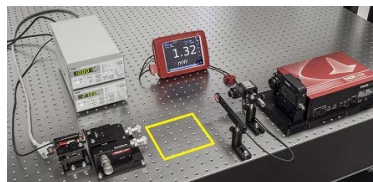
- Mounting Laser Diodes
- Driving a Laser Diode
- Selecting a Collimating Lens
- Aspheric Lenses
- Spatial Filters

The characteristics of the beams output by the different circularization systems were evaluated by making measurements using a power meter, a wavefront sensor, and an M² system. In the image of the experimental setup, all of these systems are shown on the right side of the table for illustrative purposes; they were used one at a time. The power meter was used to determine how much the beam circularization system attenuated the intensity of the input laser beam. The wavefront sensor provided a way to measure the aberrations of the output beam. The M² system measurement describes the resemblance of the output beam to a Gaussian beam. Ideally, the circularization systems would not attenuate or aberrate the laser beam, and they would output a perfectly Gaussian beam.

Edge-emitting laser diodes also emit astigmatic beams, and it can be desirable to force the displaced focal points of the orthogonal beam components to overlap. Of the three circularization techniques investigated in this work, only the cylindrical lens pair can also compensate for astigmatism. The displacement between the focal spots of the orthogonal beam components were measured for each circularization technique. In the case of the cylindrical lens pair, their configuration was tuned to minimize the astigmatism in the laser beam. The astigmatism was reported as a normalized quantity.

Lab Facts

[Click for Full Lab Facts Summary](#)



[Click to Enlarge](#)

The beam circularization systems were placed in the area of the experimental setup highlighted by the yellow rectangle.



[Click to Enlarge Cylindrical Lens Pair System](#)



[Click to Enlarge Anamorphic Prism Pair System](#)

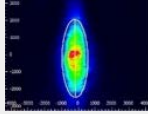
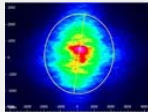
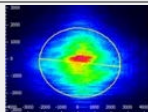
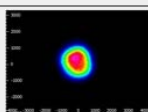


[Click to Enlarge Spatial Filter System](#)

Experimental Results

The experimental results are summarized in the following table, in which the green cells identify the best result in each category. Each circularization approach has its benefits. The best circularization technique for an application is determined by the system's requirements for beam quality, transmitted optical power, and setup constraints.

Spatial filtering significantly improved the circularity and quality of the beam, but the beam had low transmitted power. The cylindrical lens pair provided a well-circularized beam and balanced circularization and beam quality with transmitted power. In addition, the cylindrical lens pair compensated for much of the beam's astigmatism. The circularity of the beam provided by the anamorphic prism pair compared well to that of the cylindrical lens pair. The beam output from the prisms had better M^2 values and less wavefront error than the cylindrical lenses, but the transmitted power was lower.

Method	Beam Intensity Profile	Circularity ^a	M ² Values	RMS Wavefront	Transmitted Power	Normalized Astigmatism ^b
Collimated Source Output (No Circularization Technique)	 Click to Enlarge Scale in Microns	0.36	X Axis: 1.28 Y Axis: 1.63	0.17	Not Applicable	0.67
Cylindrical Lens Pair	 Click to Enlarge Scale in Microns	0.84	X Axis: 1.90 Y Axis: 1.93	0.30	91%	0.06
Anamorphic Prism Pair	 Click to Enlarge Scale in Microns	0.82	X Axis: 1.60 Y Axis: 1.46	0.16	80%	1.25
Spatial Filter	 Click to Enlarge Scale in Microns	0.93	X Axis: 1.05 Y Axis: 1.10	0.10	34%	0.36

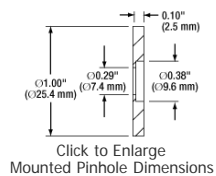
- Circularity= $d_{\text{minor}}/d_{\text{major}}$, where d_{minor} and d_{major} are minor and major diameters of fitted ellipse (1/e intensity) and Circularity = 1 indicates a perfectly circular beam.
- Normalized astigmatism is the difference in the waist positions of the two orthogonal components of the beam, divided by the Rayleigh length of the beam component with the smaller waist.

Components used in each circularization system were chosen to allow the same experimental setup be used for all experiments. This had the desired effect of allowing the results of all circularization techniques to be directly compared; however, optimizing the setup for a circularization technique could have improved its performance. The mounts used for the collimating lens and the anamorphic prism pair enabled easy manipulation and integration into this experimental system. It is possible that using smaller mounts would improve results by allowing the members of each pair to be more precisely positioned with respect to one another. In addition, using made-to-order cylindrical lenses with customized focal lengths may have improved the results of the cylindrical lens pair circularization system. All results may have been affected by the use of the beam profiler software algorithm to determine the beam radii used in the circularity calculation.

Precision Pinholes

- ▶ Pinhole Diameters from 1 μm to 1 mm
- ▶ Fabricated from Stainless Steel ($\varnothing 5 \mu\text{m}$ to $\varnothing 1 \text{ mm}$), Beryllium Copper ($\varnothing 2 \mu\text{m}$), or Black Oxide Steel ($\varnothing 1 \mu\text{m}$)
- ▶ Aluminum Housing or Black Stainless Steel Housing with 1" Outer Diameter

These mounted precision pinholes are available with pinhole diameters from 1 μm to 1 mm. Pinholes with diameters from 5 μm to 1 mm are made of 12.7 μm thick stainless steel with a black oxide coating. The 2 μm diameter pinhole is made of nickel-plated beryllium copper (12.7 μm thick), while the 1 μm diameter pinhole is made of black oxide steel (50.8 μm thick). Each pinhole is mounted in a $\varnothing 1"$, 0.10" (2.5 mm) thick aluminum disk or black stainless steel disk. Each housing is engraved with the pinhole diameter. The housing of the 1 μm pinhole is additionally engraved with the item number. The pinhole can be taken out of the housing by removing the retaining ring using a small tweezer or plier; use care as the pinhole is very thin.



Click to Enlarge
Mounted Pinhole Dimensions



Click to Enlarge
Rear of Mounted
Pinhole

Item #	Pinhole Diameter	Diameter Tolerance	Pinhole Material	Housing Material
P1H	1 μm	+0.5/-0.0 μm	Black Oxide Steel, 50.8 μm (0.002") Thick	Aluminum

P2S	2 µm	±0.5 µm	Beryllium Copper, 12.7 µm (0.0005") Thick	
P5H	5 µm	±1 µm	Stainless Steel, 12.7 µm (0.0005") Thick	Black Stainless Steel
P10H	10 µm			
P15H	15 µm	±1.5 µm		
P20H	20 µm			
P25H	25 µm			
P30H	30 µm	±2 µm		
P40H	40 µm			
P50H	50 µm	±3 µm		
P75H	75 µm			
P100H	100 µm			
P150H	150 µm	±6 µm		Aluminum
P200S	200 µm	±6 µm		
P300H	300 µm	±8 µm		
P400H	400 µm	±10 µm		Black Stainless Steel
P500H	500 µm			
P600H	600 µm			
P700H	700 µm			
P800H	800 µm			
P900H	900 µm			
P1000H	1000 µm			

Part Number	Description	Price	Availability
P1H	NEW! Ø1" Mounted Precision Pinhole, 1 +0.5/-0 µm Pinhole Diameter	\$125.00	Lead Time
P2S	Customer Inspired! Ø1" Mounted Precision Pinhole, 2 ± 0.5 µm Pinhole Diameter	\$121.00	Today
P5H	NEW! Ø1" Mounted Precision Pinhole, 5 ± 1 µm Pinhole Diameter	\$74.50	Today
P10H	NEW! Ø1" Mounted Precision Pinhole, 10 ± 1 µm Pinhole Diameter	\$74.50	Today
P15H	NEW! Ø1" Mounted Precision Pinhole, 15 ± 1.5 µm Pinhole Diameter	\$74.50	Today
P20H	NEW! Ø1" Mounted Precision Pinhole, 20 ± 2 µm Pinhole Diameter	\$67.50	Lead Time
P25H	NEW! Ø1" Mounted Precision Pinhole, 25 ± 2 µm Pinhole Diameter	\$67.50	Today
P30H	NEW! Ø1" Mounted Precision Pinhole, 30 ± 2 µm Pinhole Diameter	\$67.50	Today
P40H	NEW! Ø1" Mounted Precision Pinhole, 40 ± 3 µm Pinhole Diameter	\$67.50	Today
P50H	NEW! Ø1" Mounted Precision Pinhole, 50 ± 3 µm Pinhole Diameter	\$67.50	Today
P75H	NEW! Ø1" Mounted Precision Pinhole, 75 ± 3 µm Pinhole Diameter	\$67.50	Lead Time
P100H	NEW! Ø1" Mounted Precision Pinhole, 100 ± 4 µm Pinhole Diameter	\$67.50	Today
P150H	NEW! Ø1" Mounted Precision Pinhole, 150 ± 6 µm Pinhole Diameter	\$67.50	Today
P200S	Customer Inspired! Ø1" Mounted Precision Pinhole, 200 ± 6 µm Pinhole Diameter	\$67.50	Lead Time
P300H	NEW! Ø1" Mounted Precision Pinhole, 300 ± 8 µm Pinhole Diameter	\$67.50	Today
P400H	NEW! Ø1" Mounted Precision Pinhole, 400 ± 10 µm Pinhole Diameter	\$67.50	3-5 Days
P500H	NEW! Ø1" Mounted Precision Pinhole, 500 ± 10 µm Pinhole Diameter	\$67.50	Today
P600H	NEW! Ø1" Mounted Precision Pinhole, 600 ± 10 µm Pinhole Diameter	\$67.50	Today
P700H	NEW! Ø1" Mounted Precision Pinhole, 700 ± 10 µm Pinhole Diameter	\$67.50	Today
P800H	NEW! Ø1" Mounted Precision Pinhole, 800 ± 10 µm Pinhole Diameter	\$67.50	3-5 Days
P900H	NEW! Ø1" Mounted Precision Pinhole, 900 ± 10 µm Pinhole Diameter	\$67.50	Today
P1000H	NEW! Ø1" Mounted Precision Pinhole, 1000 ± 10 µm Pinhole Diameter	\$67.50	Today

High-Power Precision Pinholes

- ▶ Precision Copper Pinholes:
 - ▶ Gold-Plated One Side
 - ▶ Flat Poly Black (98% Emissivity) on the Reverse Side
- ▶ 25 µm Thickness at Aperture
- ▶ Black Stainless Steel Housing with 1" Outer Diameter
- ▶ High Damage Threshold:



Click to Enlarge
Gold Surface of High-Power Mounted Pinhole

Item #	Pinhole Diameter	Diameter Tolerance	Pinhole Thickness	Housing Material
P10CH	10 µm	±1 µm	25 µm (0.001")	Black Stainless Steel
P25CH	25 µm	±2 µm		
P50CH	50 µm	±3 µm		

- 5 x 10⁵ W/mm², 75 ns Pulse @ 700 nm
- 1 x 10⁶ W/mm², 10 ns Pulse @ 700 nm
- 10 W/mm², CW @ 10.6 µm

These high-power precision pinholes are designed to withstand high power densities and should be used with the beam incident on the gold-plated side. We recommend aligning the pinhole at low power, increasing the laser to full power after ensuring good throughput.

Each pinhole is mounted in a Ø1", 0.10" (2.5 mm) thick black stainless steel disk that is engraved with the diameter of the pinhole. The pinhole can be taken out of

the housing by removing the retaining ring using a small tweezer or plier; use care as the pinhole is very thin.

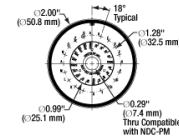
Part Number	Description	Price	Availability
P10CH	NEW! Ø1" Mounted High-Power Precision Pinhole, 10 ± 1 µm Pinhole Diameter	\$122.00	Today
P25CH	NEW! Ø1" Mounted High-Power Precision Pinhole, 25 ± 2 µm Pinhole Diameter	\$122.00	Today
P50CH	NEW! Ø1" Mounted High-Power Precision Pinhole, 50 ± 3 µm Pinhole Diameter	\$122.00	Today

Pinhole Wheels (16 Pinholes)

- ▶ Wheels with 16 Lithographically-Etched Pinholes (Transparent Glass)
- ▶ Pinhole Sizes from Ø25 µm to Ø2 mm
- ▶ Ø2.00" (Ø50.8 mm), 0.02" (0.5 mm) Thick Chrome-Plated Fused Silica Substrate
- ▶ Available Unmounted or with NDC-PM for Mounting (See the *Assembly* Tab for Mounting Instructions)
- ▶ AR Coated for 350 - 700 nm on Both Sides, $R_{avg} < 0.5\%$



Click to Enlarge
[APPLIST]
[APPLIST]
PHWM16 Post Mounted to
PY005 5-Axis Translation
Stage for Alignment



Click to Enlarge
The inner ring is engraved
with pinhole sizes and
alignment marks.

These pinhole wheels are chrome-plated fused silica disks with 16 lithographically etched pinholes ranging from Ø25 µm to Ø2 mm. The radially positioned pinholes enable a user to test multiple pinhole sizes within their experiment and requires only minor alignment after each rotation. Two versions are available; the unmounted wheel itself (Item # PWH16) and a version that includes an assembly for post mounting (Item # PHWM16). Additionally, Thorlabs offers a motorized pinhole wheel (Item # MPH16) designed for confocal microscopy systems.

The Ø2.00" (Ø50.8 mm), 0.02" (0.5 mm) thick disks are manufactured using photolithography; therefore, the pinholes are formed from the transparent substrate material where the chrome plating has been chemically etched away. Both sides of the wheel are AR coated for 350 - 700 nm ($R_{avg} < 0.5\%$) to increase transmission through the wheel. The Delrin™ inner disc is engraved with the pinhole sizes and alignment marks that point towards the pinholes. See the table below and the diagram above to the right for pinhole sizes and positions.

The NDC-PM Post Mount Assembly (included with the PHWM16) allows the pinhole wheel to be mounted to a rotating axle that can be threaded onto any Thorlabs Ø1/2" Post. Adapters for both 8-32 and M4 mounting holes are included with each mount. The shaft of the mount can be locked by tightening the side-located setscrew using the included 0.05" (1.3 mm) hex key. Assembly of the mount is required and instructions for mounting the pinhole wheel to a NDC-PM are provided in the *Assembly* tab.

Position ^a	Pinhole Size	Position	Pinhole Size	Position	Pinhole Size
A	Ø100 µm with Ø50 µm Obstruction	G	Ø50 µm	M	Ø125 µm
B	Ø25 µm	H	Ø60 µm	N	Ø200 µm
C	Ø30 µm	I	Ø70 µm	O	Ø300 µm
D	Ø35 µm	J	Ø80 µm	P	Ø1000 µm
E	Ø40 µm	K	Ø90 µm	Q	Ø2000 µm
F	Ø45 µm	L	Ø100 µm	-	-

Part Number	Description	Price	Availability
PHW16	16-Position Pinhole Wheel, Ø25 µm to Ø2 mm, Unmounted	\$416.16	Today
PHWM16	16-Position Pinhole Wheel, Ø25 µm to Ø2 mm, Mounted	\$468.18	3-5 Days